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# The Crystal Structure of Nb<sub>2</sub>Se\*

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The crystal structure of Nb<sub>2</sub>Se has been determined from single-crystal X-ray diffraction data. The compound crystallizes in the monoclinic space group C2/m with unit-cell dimensions: a=13.992, b=3.422, c=9.283 Å,  $\beta=91.76^{\circ}$ . The unit cell contains eight formula units. Refinement of the structure was performed by least-squares treatment of positional parameters and anisotropic temperature factor coefficients with three-dimensional data. All niobium atoms have distorted cubic coordination with additional near neighbors adjacent to the cube faces. The selenium atoms are bonded to between five and six niobium atoms. There is no selenium-selenium bonding in the compound. The similarities and differences between Nb<sub>2</sub>Se and other metal-rich chalcogen compounds such as Nb<sub>21</sub>S<sub>8</sub>, Ti<sub>5</sub>Te<sub>4</sub> and Ti<sub>2</sub>S, are discussed.

# Introduction

The monoselenide region of the niobium-selenium system has been investigated by Brixner (1962), by Revolinsky, Brown, Beerntsen & Armitage (1965), by Schäfer & Fuhr (1965), and by Selte, Bjerkelund & Kjekshus (1966). The region Se/Nb < 1, however, has received relatively little attention. Selte & Kjekshus (1963) reported a phase, Nb<sub>5</sub>Se<sub>4</sub>, isostructural with Ti<sub>5</sub>Te<sub>4</sub> (Grønvold, Kjekshus & Raaum, 1961). They indexed all lines on the powder pattern of Nb<sub>5</sub>Se<sub>4</sub> on a tetragonal basis with parameters: a=9.871, c=3.4529 Å. No other compounds have been reported below the ratio Se/Nb=0.8.

The application of high temperature techniques has been successful in the preparation of many metal-rich chalconides (*cf.* Franzen, Smeggil & Conard, 1967). The crystal structures of these compounds have been consistent with an interpretation of the bonding in mono- and lower chalconides of transition metals (Franzen, 1966). Comparisons of the metal-rich compounds of the Group IV transition metals with those of Group V have suggested that the roles of the chalcogen in bonding in these groups of compounds are somewhat different. For example, compare  $Ta_2S$  (Franzen & Smeggil, 1969) with the Ti<sub>2</sub>S structure type (Owens, Conard & Franzen, 1967).

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The purpose of the present work was to investigate the metal-rich region of the Nb–Se system, with the use of high temperature preparative techniques, to ascertain whether a metal-rich selenide of niobium could be synthesized, and to determine whether its structure was in accord with the above mentioned ideas concerning bonding in metal-rich chalconides and phosphides.

### Experimental

The method of high temperature preparation has been described previously (Franzen, Smeggil & Conard, 1967). The Nb<sub>2</sub>Se sample was prepared at a final annealing temperature of 1515°C. The niobium metal was obtained from E. I. Du Pont de Nemours & Co. The 99.999% selenium was obtained from Alfa Inorganics, Inc.

X-ray investigation of the sample was performed with Guinier powder and Weissenberg single-crystal techniques. For the Weissenberg exposures the crystal

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was rotated around the unique axis. The single-crystal X-ray diffraction data used in the structure determination were collected with a Hilger-Watts four circle automatic diffractometer coupled to an SDS 910 computer (Dahm, Benson, Nimrod, Fitzwater & Jacobson, 1967). Mo Ka radiation ( $\lambda = 0.7107$  Å), filtered by Zr, was used to measure the integrated intensities of 1563 reflections at  $2\theta < 80^{\circ}$ . Small corrections (less than 5%) for the non-linearity of the detector were applied to the strongest observed intensities. With the use of parameters obtained from the single-crystal diffractometer, accurate lattice parameters were determined by a leastsquares treatment of the  $2\theta$  values from Guinier powder diffraction patterns obtained at 25°C. Cu Ka radiation ( $\lambda = 1.54050$  Å), filtered by Ni, was used to obtain Guinier patterns. KCl was employed as an internal standard, a = 6.2930 Å at  $25^{\circ}$ C (Hambling, 1953). Table 1 contains the  $\sin^2\theta$  observed,  $\sin^2\theta$  calculated, and the relative intensity of each reflection from the Guinier powder pattern.

Table 1. Guinier X-ray data for Nb<sub>2</sub>Se Monoclinic, a=13.992, b=3.422, c=9.283 Å,  $\beta=91.76^{\circ}$ .

hkl	$\sin^2 \theta(\text{obs})$	$\sin^2 \theta$ (calc)	$I/I_0$	
001	0.00678	0.00689	10	
200	0.01208	0.01214	10	
201	0.01968	0.01959	5	
310	0.07802	0.07798	10	
112	0.08171	0.08181	5	
311	0.08577	0.08568	20	
312	0.10338	0.10385	100	
40 <del>3</del>	0.10712	0.10718	80	
312	0.10/12	0.10723	00	
600	0.10057	0.10918	100	
004	j 0.10937	0.11028	100	
60T	0.11441	0.11441	20	
510	0.12639	0·12651	1	
51T	0.13213	0.13201	5	
313	0.13684	0.13743	90	
020	0.20252	0.20266	70	
80 <u>2</u>	0·21716	0.21716	5	
620	0.31189	0.31189	50	
62T	0.31699	0.31707	20	
223	0.38455	0.38430	30	

Lorentz and polarization corrections were applied. From photographs ( $\times 250$ ) of the crystal, it was estimated that the shape of the crystal was a parallelepiped with one corner missing, and its size was  $88 \times 180 \times 60\mu$ . Absorption corrections were made by the use of a program by Busing & Levy (1957) adapted to the IBM 360/50 computer. Transmission factors ranged between 0.14 and 0.37. No extinction corrections were made.

Composition analysis of the sample characterized as Nb<sub>2</sub>Se was performed by igniting the sample to Nb<sub>2</sub>O<sub>5</sub>; the Se/Nb ratio was  $0.44 \pm 0.01$ . The low value for this ratio, compared with the structurally determined value of 0.5, was probably a result of the presence of a small amount of Nb metal in the Nb<sub>2</sub>Se sample, as evidenced by several faint lines observed in the powder pattern. In powder photographs of other samples lines from Nb<sub>5</sub>Se<sub>4</sub> were observed together with lines from Nb<sub>2</sub>Se. The density of Nb<sub>2</sub>Se was measured by pycnometric techniques.

### Crystal data

Nb<sub>2</sub>Se, F. W. 264.77. Monoclinic,  $a=13.992\pm0.006$ ,  $b=3.422\pm0.001$ ,  $c=9.283\pm0.007$  Å,  $\beta=91.76\pm0.06^{\circ}$ , V=446.8 Å<sup>3</sup>,  $D_m=7.8$  g.cm<sup>-3</sup>±0.1, Z=8,  $D_x=7.88$ g.cm<sup>-3</sup>. Mo K $\alpha$  ( $\lambda=0.7107$  Å)  $\mu=275.2$  cm<sup>-1</sup>. Absent reflections, hkl when h+k odd. Space group C2 (No.5), Cm(No. 8), or C2/m(No. 12). An intensity distribution analysis indicated C2/m, and this space group was confirmed by refinement of the structure.

#### Determination and refinement of the structure

A three-dimensional Patterson synthesis was performed with the use of all data; maxima were observed only at sections y=0 and  $y=\frac{1}{2}$ . The layering in the short axis direction (b=3.422 Å) was consistent with atoms occupying fourfold positions (i) x, 0, z of C2/m. Accordingly, superposition methods were performed using the y=0 section. None of the trial structures obtained in this way refined satisfactorily with the use of the 166 reflections with the largest U values (cf. below). As was discovered later, the failure of the trial structure to refine was not the result of an incorrect model, but rather the result of not including in the refinement calculation a set of reflections of relatively low intensity.

Unitary structure factors were obtained from the observed structure factors by a graphical method (Woolfson, 1961). The unitary structure factors were found to depend on k only so far as k was even or odd, *i.e.* U(h0l) = U(h2l) and U(h1l) = U(h3l), in agreement with the proposed layering in the y direction; accord-

#### Table 2. Refined atomic parameters for Nb<sub>2</sub>Se

The positional parameters and isotropic temperature factor coefficients were refined using 578 reflections of the zero and first levels. The anisotropic thermal parameters were refined using all data. All atoms occupy fourfold positions (i) x, 0, z of space group C2/m,  $\beta_{12} = \beta_{23} = 0$ .

			В				
	x	Ζ	(Ų)	$\beta_{11}$	$\beta_{22}$	β <sub>33</sub>	$\beta_{13}$
Nb(1)	0.2155 (1)	0.1251 (2)	0.43 (3)	0.00049 (4)	0.0080 (8)	0.00128 (9)	0.000?2 (4)
Nb(2)	0.0839 (1)	0.3988 (2)	0.39 (3)	0.00048 (4)	0.0064 (8)	0.00115 (9)	0.00028 (4)
Nb(3)	0.4127(2)	0.3491 (2)	0.58 (4)	0.00063 (4)	0.0076 (8)	0.0020 (1)	0.00051 (4)
Nb(4)	0.2568 (1)	0.6228 (2)	0.32(3)	0.00039 (4)	0.0058 (8)	0.00098 (9)	0.00015 (4)
Se(5)	0.1410 (2)	0.8627 (2)	0.44 (4)	0.00058 (5)	0.0065 (8)	0.0015 (1)	0.00023 (5)
Se(6)	0.4253 (2)	0.8157 (2)	0.54 (4)	0.00060 (5)	0.0077 (8)	0.0018 (1)	0.00041 (5)

ingly, only zero and first level data were used to assign phases. A set of eleven signs (Woolfson, 1961) was determined with a Harker-Kasper inequality (Harker & Kasper, 1948). The triple product sign relationship proposed by Zachariasen (1952),

$$s(\mathbf{h}) \simeq s\{\sum_{\mathbf{h}'} s(\mathbf{h}') \ s(\mathbf{h} - \mathbf{h}')\}, \qquad (1)$$

was used to extend the assignment of signs. In the case of an end-centered monoclinic crystal, the signs of two structure factors can be arbitrarily assigned (Haupt-

Table 3. Observed and calculated structure factors for Nb<sub>2</sub>Se

$ \begin{array}{c} H \\ \mathsf$		H 212213 1-1019121222473 -0000000 1201222 1-10191222 1-101912121 1-1019222 1-1019122 1-10191222 1-10191222 1-202000000000000000000000000000000	±±±±±±±±±±±±±±±±±±±±±±±±±±±±±±±±±±±±	Forststein substant for state in the st

mann & Karle, 1953);  $U_{3,1,2}$  and  $U_{\overline{6},0,1}$  were taken to be positive.

After 25 signs had been determined with the use of the eleven signs, four electron density maps (y=0)were Fourier synthesized. These four maps resulted from assignment of signs to the two symbols used in the symbolic addition procedure. All maps had peaks in basically the same relative positions; one was selected as a trial structure on the basis of its similarity with a map obtained by the superposition method. Structure factors were then assigned signs and added to those originally used in groups of 10–20, in repeated attempts to refine the trial structure discussed above, and to obtain new electron density sections at y=0. Finally, after accepting 149 reflections, the structure refined to a reliability index

$$R \equiv \frac{\Sigma \left| |F_o| - |F_c| \right|}{\Sigma \left| F_o \right|} = 0.14.$$

The positional parameters and isotropic thermal parameters of the six independent atoms of the structure were refined by the computer program ORFLS (Busing, Martin & Levy, 1962) with 578 reflections of the zero and first levels, and the 020 reflection. The h0l and h1l data were accepted as 'observed' provided that  $\sigma(I)/I < 0.4$  $[\sigma(I)$  obtained as discussed below]. The atomic scattering factors used were those given by Hanson, Herman, Lea & Skillman (1964); both real and imaginary components of anomalous dispersion were included for Nb, -2.1 and 0.9, respectively, and for Se, -0.1 and 2.4, respectively (values given are for sin  $\theta/\lambda = 0.0$ , International Tables for X-ray Crystallography, 1962). The function minimized in the least-squares procedure was  $\Sigma w(|F_o| - |F_c|)^2$ , where  $F_o$  and  $F_c$  are observed and calculated structure factors and  $w = 1/\sigma^2(F_o)$ . The  $\sigma(F_o)$ was initially obtained from the following relationship

$$\sigma^{2}(F_{o}) = \frac{A^{*}\mathrm{Lp}}{4} \left( \frac{I_{\mathrm{B}} + I_{T}}{I_{T} - I_{\mathrm{B}}} \right) + \frac{\mathrm{Lp}(I_{T} - I_{\mathrm{B}})}{4A^{*}} \sigma^{2}(A^{*})$$
$$+ \frac{(I_{T} - I_{\mathrm{B}})A^{*}}{4\mathrm{Lp}} \sigma^{2}(\mathrm{Lp}), \qquad (2)$$

where  $A^* =$  absorption factor, Lp=Lorentz-polarization correction factor,  $I_T =$  total counts,  $I_B =$  background counts,  $\sigma^2(A^*) = 0.0025$ , and  $\sigma^2(2\theta) = 0.0004$ . The last quantity was used to calculate  $\sigma^2(Lp)$ . The weighting scheme was checked by plotting  $w(F_o - F_c)^2$ versus  $F_o^2$  and fitting the data to a straight line by a least-squares technique. The initial weights, as calculated from (2), gave a line with slope 2.3. The weights were adjusted so that the straight line of least-squares fit had a slope of 0.01.

An unweighted R value of 0.105 was obtained for the 578 reflections after several cycles; for the last cycle all parameter shifts were < 1% of the estimated standard deviations of the parameters. The final values of

the parameters and their standard deviations are given in Table 2. The standard deviation of an observation of unit weight,

$$[\Sigma w \Delta^2/m - n]^{1/2}, \qquad (3)$$

was 1.03. In expression (3), w is the revised weight of an observation,  $\Delta$  is  $(F_o - F_c)$ , m is the number of observations (578), and n is the number of variables (19). The positional parameters obtained from the refinement described above were used to calculate all structure factors in the zero and first levels, including those observations which had been taken as unobserved, as discussed above. The list of observed and calculated

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structure factors are tabulated in Table 3. The reliability index for these data was 0.111.

Refinement of the structure was also performed with anisotropic thermal parameters by use of all 1563 data (maximum k=6). The anisotropic components of the thermal parameters are given in Table 2. The positional parameters obtained in this refinement were the same, within the standard deviations, as those listed in Table 2. In this case the *R* value was 0.100, and the 'goodness of fit', as defined by (3), was 0.94.\* Refinements were

\* A total listing of  $F_{obs}$  and  $F_{calc}$  for the anisotropic refinement will be furnished upon request.

<b>D</b> (		No. of		
Reference atom	Neighbor	neighbors	Distance	
Nb(1)	Se(5)	1	2.618	
	Se(S)	2	2.642	
	Nb(4)	2	2.082	
	Nb(1)	2	3.074	
	Nb(2)	1	3.197	
	Nb(3)	ĩ	3.404	
Nb(2)	Se(6)	2	2.630	
	Nb(4)	2	2.826	
	Nb(3)	2	2.904	
	ND(3)	2	2.907	
	Nb(2)	1	3.138	
	Nb(1)	1	3.197	
Nb(3)	Se(5)	2	2.701	
	Se(6)	1	2.782	
	Nb(2)	2	2.904	
	Nb(4)	2	2.947	
	ND(2) Nb(1)	2	3,404	
	Nb(4)	· 1	3.411	
Nb(4)	Se(5)	1	2.805	
	Nb(2)	2	2.826	
	Nb(4)	2	2.857	
	Se(6)	1	2.911	
	Nb(1)	$\frac{2}{2}$	2.925	
	Nb(2)	1	3.138	
	Nb(3)	î	3.411	
Se(5)	Nb(1)	1	<b>2</b> ·618	
	Nb(1)	2	2.642	
	Nb(3) Nb(4)	2	2·701 2·805	
Sa(6)	NL(2)	1	2.630	
Se(6)	Nb(2)	2	2.630	
	Nb(3)	1	2.782	
	Nb(4)	ī	2.911	
Angle I	Degrees	1	Angle	Degrees
(1)-Se(5)-Nb(1)	71.51	Nb(3)-5	Se(6)-Nb(2)	66•41
(1)-Se(5)-Nb(1)	80.96	Nb(2)-	Se(6)–Nb(2)	81.39
(3)-Se(5)-Nb(3)	78.84	Nb(1)-	Se(6)-Nb(1)	79.49
(1) - Se(5) - Nb(4)	121.36	Nb(3)-	Se(6) - Nb(4)	108.53
(1) - Se(3) - IND(4) (1) - Se(5) - Nib(3)	04.83 140.53	ND(2)-	Se(0) - IND(4) Se(6) - Nb(1)	01.07
(1) - Se(5) - Nb(3)	79.16	Nb(2) =	Se(6) - Nb(1)	73.99
(3) - Se(5) - Nb(4)	64.70	Nb(1)-	Se(6) - Nb(4)	62.86

Table 4. Interatomic distances in Nb<sub>2</sub>Se ( $\pm 0.003$  Å) and angles around Se atoms

also carried out in the acentric space groups. In each case the hypothesis that the acentric structure was correct rather than the centric one could be rejected according to Hamilton's (1965) criterion at the 99.5% confidence level.

For the purpose of checking the correctness of the refinement, a difference electron density synthesis was performed with 578 reflections of the zero and first levels. The synthesis had maxima, corresponding to about 2 electrons at the y=0 section; these maxima were all at atomic positions. Interatomic distances and angles, obtained by the use of the *Thermal Ellipsoid Plot* computer program written by Johnson (1965), are listed in Table 4. Johnson's program was used with Slater's (1964) atomic radii to draw the stereoscopic view of the structure along the *b* axis. The illustration is presented in Fig.1.

#### Description and discussion of the structure

The four independent Nb atoms have similar coordination polyhedra, namely, a distorted cubic arrangement of eight atoms with additional bonded atoms adjacent to the faces of the cube. For example, as shown in Fig.2, the most distorted cubic arrangement is around the Nb(1) atom which is coordinated by 2 Nb(1), 2 Nb(4), 2 Se(5), and 2 Se(6), forming the distorted cube, and Nb(2), Nb(3), and Se(5) are located adjacent to the faces of the cube. The least distorted cubic arrangement is found around Nb(4).

The coordination polyhedra of the two independent Se atoms are difficult to describe. One can see, however, by comparing angles and distances in Table 4 that the arrangement of Nb atoms around each Se is similar, namely, that each Se atom is bonded to six Nb atoms. In addition, it is noted that the Nb atoms are not arranged so as to spatially surround the Se atoms; rather, in the vicinity of each Se there are relatively large regions which are unoccupied by either Nb or Se atoms. Referring to Fig. 2, these 'holes' occur on the twofold axes.

The Nb<sub>2</sub>Se structure-type is new, but the metal and non-metal environments in Nb<sub>2</sub>Se are similar to those

in Nb<sub>5</sub>Se<sub>4</sub>, which has the Ti<sub>5</sub>Te<sub>4</sub> type structure (Selte & Kjekshus, 1963). The coordination polyhedra of the non-metal in Nb<sub>2</sub>Se and Nb<sub>5</sub>Se<sub>4</sub> are related to those observed in Nb<sub>3</sub>Se<sub>4</sub> (Selte & Kjekshus, 1964), perhaps suggesting a basis for understanding the relationship of metal-rich chalconides to mono- and higher chalconides.

The tendency for metal atoms to form body-centered coordination polyhedra has been correlated with the promotion energies from the ground electronic states to the  $sd^{n-1}$  configurations (*n* denotes the number of valence electrons) in the gaseous atoms (Brewer, 1963). The metal coordination in metal-rich compounds has recently been discussed from this point of view (Franzen, Smeggil & Conard, 1967). The Nb<sub>2</sub>Se structure type provides further corroboration of this correlation, for elemental Nb crystallizes in the b.c.c.structure type, and the Nb atoms in Nb<sub>2</sub>Se are surrounded by eight atoms, both Se and Nb, in distorted cubes.

The Nb<sub>2</sub>Se structure can be compared with the  $Nb_{21}S_8$  structure (Franzen, Beineke & Conard, 1968). Fig.2 shows the Nb<sub>2</sub>Se structure projected along the short axis (3.42 Å). Fig. 3 shows the  $Nb_{21}S_8$  structure also projected along its short axis (3.36 Å). There is a mirror plane parallel to the plane of the Figure in each structure. The structural similarity between Nb<sub>2</sub>Se and  $Nb_{21}S_8$  is outlined by exactly equivalent bold-lined parallelograms in each Figure. In Nb<sub>2</sub>Se the bold parallelogram is almost an asymmetric unit of the structure, whereas in  $Nb_{21}S_8$  the same parallelogram is only a portion of an asymmetric unit. The projected atom positions in each bold-lined parallelogram are approximately the same with the exception that in Nb<sub>21</sub>S<sub>8</sub> there are two projected atom positions which are not present in Nb<sub>2</sub>Se, namely, the Nb atom at z=0 in the upper right-hand corner and the Nb atom at  $z = \frac{1}{2}$  in the lower right-hand corner of the bold-lined parallelogram. It is seen, however, that these two additional projected positions in  $Nb_{21}S_8$  occur in the 'hole' regions of  $Nb_2Se$ . Nonetheless, Nb<sub>2</sub>Se is seen to have a basic structural similarity with Nb<sub>21</sub>S<sub>8</sub>, and thus with a number of lower chalconides (and, it should be mentioned, phosphides) as well.



Fig. 1. Stereoscopic illustration of the Nb<sub>2</sub>Se structure viewed along the b axis. The box indicates a unit cell.

Nb<sub>2</sub>Se does exhibit, however, some significant differences from other metal-rich transition metal compounds. For example, the  $Nb_{21}S_8$  structure does not contain regions with non-bonded sulfur arrangements. On the other hand, while, as is generally true of metalrich chalconides, there are no selenium-selenium bonds in the Nb<sub>2</sub>Se structure (the closest Se-Se distance is 3.43 Å), the Se atoms are, nonetheless, arranged in nonbonded tetrahedra. These tetrahedra share edges so as to form chains of tetrahedra in the b direction. The non-bonded Se atom arrangements in Nb<sub>2</sub>Se differentiate this structure from the structures of the Group IV transition metal subchalconides and subphosphides, for in all of the Group IV compounds a high degree of coordination (six to nine) of the non-metal by metal atoms is observed. This feature of Nb<sub>2</sub>Se is shared by Ta<sub>2</sub>S (Franzen & Smeggil, 1969), for in Ta<sub>2</sub>S the sulfur atoms are similarly located in a non-bonded fashion in regions of the structure.

On the basis of these and similar observations, the following classification can be made of the metal-rich chalconides of transition metals: Class I refers to group IV transition metal compounds with chalcogen/metal (C/M)=0.5, which form with high chalcogen coordination numbers (six to nine); to Class II belong the Group V compounds with C/M < 0.5 and high chalcogen coordination numbers; and to Class III belong other Group V compounds which have C/M=0.5 and lower chalcogen coordination numbers of Class III compounds exhibit regions containing non-bonded chalcogen atoms and regions with extensive metal-metal bonding.

# Table 5. Known metal-rich chalconide compounds, listed according to their class (defined in text)

Average chalcogen coordination numbers (C.N.) and calculated chalcogen valences are given.

	Class	Average C.N.	Average calculated valence
Ti <sub>2</sub> S	Ι	7.7	4.5
$Ti_2Se^{(a)}$	I	7.7	<u> </u>
$Zr_2S^{(a)}$	I	77	
$Zr_2Se^{(b)}$	ĩ	·•7	4.3
$Hf_2S^{(c)}$	1	6	3.4
$Hf_2Se^{(a)}$	j j	6	
$\alpha - \bar{V}_3 S^{(d)}$	II	8	5.0
$\beta - V_3 S^{(d)}$	II	8	5.0
Nb <sub>2</sub> S <sub>3</sub>	II	7.5	4.1
Nb <sub>2</sub> Se	III	6	3.0
Ta <sub>2</sub> S	III	5	3.0

References not given in text:

(a) Franzen, Smeggil & Conard (1967).

(b) Franzen & Norrby (1968).

(c) Franzen & Graham (1966).
 (d) Pedersen & Grønvold (1959).

Table 5 lists all known subchalconides of the Group IV and V transition metals, together with their classes (as defined above), their average chalcogen coordination numbers, and their average chalcogen valences, as given by the Pauling empirical relationship (Pauling, 1960). While these valences are probably not absolutely correct, the trends in the calculated valences clearly are significant in this case.

The coordination numbers and valences of the chalcogens in the various classes strongly suggest an elec-



Fig. 2. A unit cell of Nb<sub>2</sub>Se projected on (010). Bold-lined parallelogram contains structural similarities to Nb<sub>21</sub>S<sub>8</sub>.

tron concentration interpretation of the differences between the classes. All Group IV subchalconides belong to Class I, as would be expected from this viewpoint, since all have the same valence electron concentration. On the other hand, the valence electron concentration is increased if the stoichiometry is unchanged but a Group V element replaces a Group IV element. Since no Group V subchalconides are isostructural with Group IV subchalconides, this replacement appears to result in compounds unstable either with respect to disproportionation into a still lower chalconide (forming a Class II compound) and a higher chalconide, or with respect to a structure type in which the chalcogen contributes fewer electrons to the conduction band (forming a Class III compound).

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Fig. 3. Approximately one and one-half unit cells of  $Nb_{21}S_8$  projected on (001). All sulfur atoms are coordinated to six niobium atoms at the corners of trigonal prisms, shown by the dotted lines, and to additional niobium atoms adjacent to rectangular faces of the prisms. The bold-lined parallelogram contains structural similarities to  $Nb_2Se$ .

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# The Crystal Structure of Ta<sub>2</sub>S\*

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The crystal structure of ditantalum sulfide, Ta<sub>2</sub>S, has been determined from single-crystal X-ray diffraction data. The compound crystallizes in the orthorhombic system, space group *Pbcm*, with unitcell dimensions: a=7.379, b=5.574, and c=15.19 Å. The unit cell contains 12 formula units. Refinement of the structure was performed by a least-squares treatment of the positional parameters and the temperature factor coefficients. The Ta<sub>2</sub>S structure is fundamentally different from other known metalrich chalconide structures. The structure can be viewed as chains of body-centered pentagonal antiprisms of tantalum atoms sharing faces in one direction and interconnected *via* sulfur atoms in the other two directions.

### Introduction

The crystal structures of the sulfides of tantalum have been reviewed by Jellinek (1963). It was concluded that an early claim of the preparation of a subsulfide of tantalum by Biltz & Kocher (1938) was incorrectly based on powder patterns of a mixture of phases. Preparations at relatively low temperatures (less than about 1000°C) have generally failed to produce metalrich sulfides such as have been prepared by high-temperature techniques in recent years in the Hf-S system (Franzen & Graham, 1966), in the Ti-S system (Owens, Conard & Franzen, 1967), and in the Nb-S system (Franzen, Beineke & Conard, 1968). The work reported here was performed with the purpose of extending the application of high-temperature preparation techniques to the Ta-S system in order to prepare and determine the structure of a lower sulfide of tantalum.

#### Experimental

The ditantalum sulfide samples were prepared in a manner analogous to that previously described (Owens, Conard & Franzen, 1967). The sulfur obtained from the American Smelting and Refining Co. was 99.999% pure. The tantalum obtained from the National Research Corporation was 99.99% pure. The final annealing temperature in the case of the Ta<sub>2</sub>S preparation was 1600°C. Single-crystal X-ray diffraction data were initially collected with a General Electric spectrogoniometer and nickel filtered Cu Ka radiation. Data were collected for 767 reflections in the range  $0^{\circ} < 2\theta \leq$ 160°. A second set of data was collected using a Hilger-Watts diffractometer coupled with an SDS(910-IBM) (1401) computer configuration described elsewhere (Dahm, Benson, Nimrod, Fitzwater & Jacobson, 1967). Mo  $K\alpha$  radiation was used with a Zr filter to collect data for 1104 reflections in the range  $0^{\circ} < 2\theta \le 60^{\circ}$ . Lorentz and polarization corrections were applied to both sets of data. Absorption corrections, based on the crystal's approximately rectangular prismatic shape

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